

LTI Discrete-Time Systems in the Transform Domain

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Types of Transfer Functions

- The time-domain classification of an LTI digital transfer function sequence is based on the length of its impulse response:
 - Finite impulse response (FIR) transfer function
 - Infinite impulse response (IIR) transfer function
- In the case of digital transfer functions with frequencyselective frequency responses, there are two types of classifications
 - Classification based on the shape of the magnitude function $|H(e^{j\omega})|$
 - Classification based on the form of the phase function $\theta(\omega)$

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Classification Based on Magnitude Characteristics

- One common classification is based on ideal magnitude response
- A digital filter designed to pass signal components of certain frequencies without distortion should have a magnitude response equal to one at these frequencies, and zero at all other frequencies
- The range of frequencies where the frequency response takes the value of one is called the passband
- The range of frequencies where the frequency response takes the value of zero is called the **stopband**

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Ideal Filters

• Frequency responses of the four popular types of ideal digital filters with real coefficients are shown below:



- The freq. ω_c , ω_{c1} , and ω_{c2} are called the **cutoff frequencies**
- An ideal filter has a magnitude response equal to one in the passband and zero in the stopband, and has a zero phase everywhere

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Ideal Filters

• The impulse response of the ideal lowpass filter:

$$h_{LP}[n] = \frac{\sin \omega_c n}{\pi n}, \quad -\infty < n < \infty$$

- The above impulse response is not absolutely summable, and hence, the corresponding transfer function is not BIBO stable
- Also, $h_{LP}[n]$ is not causal and is of doubly infinite length
- The remaining three ideal filters are also characterized by doubly infinite, noncausal impulse responses and are not absolutely summable
- Thus, the ideal filters with the ideal "brick wall" frequency responses cannot be realized with finite dimensional LTI filter

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Ideal Filters

- To develop stable and realizable transfer functions, the ideal freq. response specs. are relaxed by including a transition **band** between the passband and the stopband
- This allows the magnitude response to decay slowly from its max. value in the passband to the zero value in the stopband

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 Moreover, the magnitude response is allowed to vary by a small amount both in the passband and the stopband



 A causal stable real-coefficient transfer function H(z) is defined as a bounded real (BR) transfer function if

 $|H(e^{j\omega})| \le 1$ for all values of ω

- Let x[n] and y[n] denote, respectively, the input and output of a digital filter characterized by a BR transfer function H(z)with $X(e^{j\omega})$ and $Y(e^{j\omega})$ denoting their DTFTs
- Then the condition $|H(e^{j\omega})| \le 1$ implies that

 $|Y(e^{j\omega})|^2 \leq |X(e^{j\omega})|^2$

Integrating the above from – π to π, and applying Parseval's relation we get

$$\sum_{n=-\infty}^{\infty} |y[n]|^2 \le \sum_{n=-\infty}^{\infty} |x[n]|^2$$

- Thus, for all finite-energy inputs, the output energy is less than or equal to the input energy implying that a digital filter characterized by a BR transfer function can be viewed as a passive structure
- If |*H*(*e^{jω}*)| = 1, then the output energy is equal to the input energy, and such a digital filter is therefore a lossless system
- A causal stable real-coefficient transfer function *H*(*z*) with |*H*(*e^{jω}*)| = 1 is thus called a **lossless bounded real (LBR)** transfer function
- The BR and LBR transfer functions are the keys to the realization of digital filters with low coefficient sensitivity (see Sec. 12.9)

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• Example – Consider the causal stable IIR transfer function

$$H(z) = \frac{K}{1 - \alpha z^{-1}}, \quad 0 < |\alpha| < 1$$

where *K* is a real constant

- Its square-magnitude function is given by $|H(e^{j\omega})|^2 = H(z)H(z^{-1})|_{z=e^{j\omega}} = \frac{K^2}{(1+\alpha^2) - 2\alpha\cos\omega}$
- The maximum value of |*H*(*e^{jω}*)|² is obtained when 2α cosω in the denominator is a maximum and the minimum value is obtained when 2α cosω is a minimum
- For α > 0, maximum value of 2α cosω is equal to 2α at ω = 0, and minimum value is -2α at ω = π

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- Thus, for $\alpha > 0$, the maximum value of $|H(e^{j\omega})|^2$ is equal to $K^2 / (1 \alpha)^2$ at $\omega = 0$ and the minimum value is equal to $K^2 / (1 + \alpha)^2$ at $\omega = \pi$
- On the other hand, for α < 0, the maximum value of 2α cosω is equal to -2α at ω = π, and the minimum value is equal to 2α at ω = 0
- Here, the maximum value of $|H(e^{j\omega})|^2$ is equal to $K^2 / (1 + \alpha)^2$ at $\omega = \pi$ and the minimum value is equal to $K^2 / (1 - \alpha)^2$ at $\omega = 0$
- Hence, the maximum value can be made equal to 1 by choosing $K = \pm (1 \alpha)$, in which case the minimum value becomes $(1 \alpha)^2/(1 + \alpha)^2$

• Hence,

$$H(z) = \frac{K}{1 - \alpha z^{-1}}, \quad 0 < |\alpha| < 1$$

is a BR function for $K = \pm(1 - \alpha)$

 Plots of the magnitude function for with values of K chosen to make H(z) a BR function are shown below



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• An IIR transfer function A(z) with unity magnitude response for all frequencies, i.e.,

 $|H(e^{j\omega})|^2 = 1$, for all ω

is called an allpass transfer function

• An *M*-th order causal real-coefficient allpass transfer function is of the form $A_M(z) = + \frac{d_M + d_{M-1}z^{-1} + \dots + d_1z^{-M+1} + z^{-M}}{d_M + d_M + d_M$

$$A_M(z) = \pm \frac{M - M - 1}{1 + d_1 z^{-1} + \dots + d_{M-1} z^{-M+1} + d_M z^{-M}}$$

• Denote the denominator polynomials of $A_M(z)$ as $D_M(z)$:

$$D_M(z) = 1 + d_1 z^{-1} + \dots + d_{M-1} z^{-M+1} + d_M z^{-M}$$

then it follows that $A_M(z)$ can be written as

$$A_M(z) = \pm \frac{z^{-M} D_M(z^{-1})}{D_M(z)}$$

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- Note from the above that if $z = re^{j\phi}$ is a pole of a real coefficient allpass transfer function, then it has a zero at $z = (1/r)e^{j\phi}$
- The numerator of a real-coefficient allpass transfer function is said to be the mirror image polynomial of the denominator, and vice versa
- We shall use the notation $\tilde{D}_M(z)$ to denote the mirror-image polynomial of a degree-*M* polynomial $D_M(z)$, i.e.,

$$\widetilde{D}_M(z) = z^{-M} D_M(z^{-1})$$

• The expression

$$A_M(z) = \pm \frac{z^{-M} D_M(z^{-1})}{D_M(z)}$$

implies that the poles and zeros of a real-coefficient allpass function exhibit mirror image symmetry in the z-plane 7-13 Original PowerPoint slides prepared by S. K. Mitra

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• To show that $|A_M(e^{j\omega})| = 1$ we observe that

$$A_M(z^{-1}) = \pm \frac{z - D_M(z)}{D_M(z^{-1})}$$

• Therefore

$$A_M(z)A_M(z^{-1}) = \frac{z^{-M}D_M(z^{-1})}{D_M(z)} \frac{z^M D_M(z)}{D_M(z^{-1})}$$

Hence

$$|A_M(e^{j\omega})|^2 = A_M(z)A_M(z^{-1})\Big|_{z=e^{j\omega}} = 1$$

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- Now, the poles of a causal stable transfer function must lie inside the unit circle in the z-plane
- Hence, all zeros of a causal stable allpass transfer function must lie outside the unit circle in a mirror-image symmetry with its poles situated inside the unit circle
- Figure below shows the principal value of the phase of the 3rd-order allpass function



• If we unwrap the phase by removing the discontinuity, we arrive at the unwrapped phase function $\theta_c(\omega)$ as follows



• The unwrapped phase function of any arbitrary causal stable allpass function is a continuous function of ω

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Properties:

- A causal stable real-coefficient allpass transfer function is a lossless bounded real (LBR) function or, equivalently, a causal stable allpass filter is a lossless structure
- 2. The magnitude function of a stable allpass function A(z) satisfies: |z| > 1

$$A(z) \bigg| \begin{cases} < 1, & \text{for } |z| > 1 \\ = 1, & \text{for } |z| = 1 \\ > 1, & \text{for } |z| < 1 \end{cases}$$

3. Let $\tau(\omega)$ denote the group delay function of an allpass filter A(z), i.e.,

$$\tau(\omega) = -\frac{d}{d\omega} [\theta_c(\omega)]$$

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- The unwrapped phase function $\theta_c(\omega)$ of a stable allpass function is a monotonically decreasing function of ω so that $\tau(\omega)$ is everywhere positive in the range $0 < \omega < \pi$
- The group delay of an *M*-th order stable real-coefficient allpass transfer function satisfies $\int_{1}^{\pi} \tau(\omega) d\omega = M\pi$
- A Simple Application:
- A simple but often used application of an allpass filter is as a delay equalizer
- Let *G*(*z*) be the transfer function of a digital filter designed to meet a prescribed magnitude response
- The nonlinear phase response of G(z) can be corrected by cascading it with an allpass filter A(z) so that the overall cascade has a constant group delay in the band of interest

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• Since $|A(e^{j\omega})| = 1$, we have

 $|A(e^{j\omega}) G(e^{j\omega})| = |G(e^{j\omega})|$

- Overall group delay is the given by the sum of the group delays of G(z) and A(z)
- Example Figure below shows the group delay of a 4th order elliptic filter with the following specifications: $\omega_p = 0.3\pi$, $\delta_p = 1 \text{ dB}$, $\delta_s = 35 \text{ dB}$
- The group delay of the original filter cascaded with an 8th order allpass filter is also shown
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Classification Based on Phase Characteristics

- A second classification of a transfer function is with respect to its phase characteristics
- In many applications, it is necessary that the digital filter designed does not distort the phase of the input signal components with frequencies in the passband
- One way to avoid any phase distortion is to make the frequency response of the filter real and nonnegative, i.e., to design the filter with a zero phase characteristic
- However, it is not possible to design a causal digital filter with a zero phase
- For non-real-time processing of real-valued input signals of finite length, zero-phase filtering can be very simply implemented by relaxing the causality requirement 7-20

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Zero-Phase Transfer Functions

One zero-phase filtering scheme is sketched below

$$x[n] \longrightarrow H(z) \longrightarrow v[n] \qquad u[n] \longrightarrow H(z) \longrightarrow w[n]$$
$$u[n] = v[-n], \qquad y[n] = w[-n]$$

- Let X(e^{jω}), V(e^{jω}), U(e^{jω}), W(e^{jω}), and Y(e^{jω})denote the DTFTs of x[n], v[n], u[n], w[n], and y[n], respectively
- Making use of the symmetry relations we arrive at the relations between various DTFTs as follows:

$$V(e^{j\omega}) = H(e^{j\omega}) X(e^{j\omega}), W(e^{j\omega}) = H(e^{j\omega}) U(e^{j\omega})$$
$$U(e^{j\omega}) = V^*(e^{j\omega}), Y(e^{j\omega}) = W^*(e^{j\omega})$$

• Combining the above equations we get

$$Y(e^{j\omega}) = W^*(e^{j\omega}) = H^*(e^{j\omega}) U^*(e^{j\omega}) = H^*(e^{j\omega}) V(e^{j\omega})$$

$$= H^*(e^{j\omega}) H(e^{j\omega}) X(e^{j\omega}) = |H(e^{j\omega})|^2 X(e^{j\omega})$$

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• The output y[n] of a linear-phase filter to an input $x[n] = Ae^{j\omega n}$ is then given by

$$y[n] = Ae^{-j\omega D} e^{j\omega n} = Ae^{j\omega(n-D)}$$

- If x_a(t) and y_a(t) represent the continuous-time signals whose sampled versions, sampled at t = nT, are x[n] and y[n] given above, then the delay between x_a(t) and y_a(t) is precisely the group delay of amount D
- If D is an integer, then y[n] is identical to x[n], but delayed by D samples
- If D is not an integer, y[n], being delayed by a fractional part, is not identical to x[n]
 - The waveform of the underlying continuous-time output is identical to the waveform of the continuous-time input and delayed *D* units of

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- If it is desired to pass input signal components in a certain frequency range undistorted in both magnitude and phase, then the transfer function should exhibit a unity magnitude response and a linear-phase response in the band of interest
- Figure below shows the frequency response if a lowpass filter with a linear-phase characteristic in the passband
- Since the signal components in the stopban are blocked, the phase response in the stopband can be of any shape



 Example – Determine the impulse response of an ideal lowpass filter with a linear phase response

$$H_{LP}(e^{j\omega}) = \begin{cases} e^{-j\omega n_o}, & 0 < |\omega| < \omega_c \\ 0, & \omega_c \le |\omega| \le \pi \end{cases}$$

- Applying the frequency-shifting property of the DTFT to the impulse response of an ideal zero-phase lowpass filter we arrive at $h_{LP}[n] = \frac{\sin \omega_c (n n_o)}{\pi (n n_o)}, \quad -\infty < n < \infty$
- As before, the above filter is noncausal and of doubly infinite length, and hence, unrealizable
- By truncating the impulse response to a finite number of terms, a realizable FIR approximation to the ideal lowpass filter can be developed

- The truncated approximation may or may not exhibit linear • phase, depending on the value of n_o chosen
- If we choose $n_o = N/2$ with N a positive integer, the truncated and shifted approximation

$$\hat{h}_{LP}[n] = \frac{\sin \omega_c (n - N/2)}{\pi (n - N/2)}, \quad 0 \le n \le N$$

will be a length *N*+1 causal linear-phase FIR filter



Zero-Phase Transfer Functions

- Because of the symmetry of the impulse response coefficients as indicated in the two figures, the frequency response of the truncated approximation can be expressed as: $\hat{H}_{LP}(e^{j\omega}) = \sum_{n=0}^{N} \hat{h}_{LP}[n]e^{-j\omega n} = e^{-j\omega N/2} \tilde{H}_{LP}(\omega)$
- Where *H˜_{LP}*(ω), called the zero-phase response or amplitude response, is a real function of ω will be a length *N*+1 causal linear-phase FIR filter

• Consider the two 1st-order transfer functions:

 $H_1(z) = \frac{z+b}{z+a}, \ H_2(z) = \frac{bz+1}{z+a}, \ |a| < 1, \ |b| < 1$

- Both transfer functions have a pole inside the unit circle at the same location and are stable
- But the zero of $H_1(z)$ is inside the unit circle at z = -b, whereas, the zero of $H_2(z)$ is at z = -1/b situated in a mirrorimage symmetry $H_1(z) = H_2(z)$



 However, both transfer functions have an identical magnitude function as

 $H_1(z)H_1(z^{-1}) = H_2(z)H_2(z^{-1})$

• The corresponding phase functions are

$$\arg[H_1(e^{j\omega})] = \tan^{-1} \frac{\sin\omega}{b + \cos\omega} - \tan^{-1} \frac{\sin\omega}{a + \cos\omega}$$

$$\arg[H_2(e^{j\omega})] = \tan^{-1} \frac{b\sin\omega}{1+b\cos\omega} - \tan^{-1} \frac{\sin\omega}{a+\cos\omega}$$

 Figure below shows the unwrapped phase responses of the two transfer functions for a = 0.8 and b = -0.5



- As shown in the figure, H₂(z) has an excess phase lag with respect to H₁(z)
- The excess phase lag property of $H_2(z)$ with respect to $H_1(z)$ can also be explained by observing that we can write

$$H_2(z) = \frac{bz+1}{z+a} = \left(\underbrace{\frac{z+b}{z+a}}_{H_1(z)} \underbrace{\frac{bz+1}{z+b}}_{A(z)}\right)$$

where A(z) = (bz + 1) / (z + b) is a stable allpass function

- The phase functions of $H_1(z)$ and $H_2(z)$ are related through $\arg[H_2(e^{j\omega})] = \arg[H_1(e^{j\omega})] + \arg[A(e^{j\omega})]$
- As the unwrapped phase function of a stable first-order allpass function is a negative function of ω , it follows from the above that $H_2(z)$ has an excess phase lag with $H_1(z)_{7-29}$

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- Generalizing the above result, let $H_m(z)$ be a causal stable transfer function with all zeros inside the unit circle and let H(z) be another causal stable transfer function satisfying $|H(e^{j\omega})| = |H_m(e^{j\omega})|$
- These two transfer functions are then related through $H(z) = H_m(z) A(z)$ where A(z) is a causal stable allpass function
- The unwrapped phase functions of $H_m(z)$ and H(z) are thus related through

 $\arg[H(e^{j\omega})] = \arg[H_m(e^{j\omega})] + \arg[A(e^{j\omega})]$

- H(z) has an excess phase lag with $H_m(z)$
- A causal stable transfer function with all zeros inside the unit circle is called a **minimum-phase transfer function**

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- A causal stable transfer function with all zeros outside the unit circle is called a **maximum-phase transfer function**
- A causal stable transfer function with zeros inside and outside the unit circle is called a mixed-phase transfer function
- Example Consider the mixed-phase transfer function $H(z) = \frac{2(1+0.3z^{-1})(0.4-z^{-1})}{(1-0.2z^{-1})(1+0.5z^{-1})}$
- We can rewrite H(z) as

$$H(z) = \left[\frac{2(1+0.3z^{-1})(1-0.4z^{-1})}{(1-0.2z^{-1})(1+0.5z^{-1})}\right] \left(\frac{0.4-z^{-1}}{1-0.4z^{-1}}\right)$$

Minimum-phase function

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Allpass function

- It is impossible to design an IIR transfer function with an exact linear-phase
- It is always possible to design an FIR transfer function with an exact linear-phase response
- We now develop the forms of the linear-phase FIR transfer function H(z) with real impulse response h[n]
- Let $H(z) = \sum_{n=0}^{N} h[n] z^{-n}$
- If *H*(*z*) is to have a linear-phase, its frequency response must be of the form

$$H(e^{j\omega}) = e^{j(c\omega+\beta)}\check{H}(\omega)$$

where *c* and β are constants, and $\check{H}(\omega)$, called the amplitude response (zero-phase response), is a real function of ω Original PowerPoint slides prepared by S. K. Mitra

- For a real impulse response, the magnitude response $|H(e^{j\omega})|$ is an even function of ω , i.e., $|H(e^{j\omega})| = |H(e^{-j\omega})|$
- Since $|H(e^{j\omega})| = |\check{H}(\omega)|$, the amplitude response is then either an even function or an odd function of ω , i.e., $\check{H}(-\omega) = \pm \check{H}(\omega)$
- The frequency response satisfies the relation

$$|H(e^{j\omega})| = |H^*(e^{-j\omega})|$$

or, equivalently, the relation

 $e^{j(c\omega+\beta)}\check{H}(\omega) = e^{-j(-c\omega+\beta)}\check{H}(-\omega)$

• If $\check{H}(\omega)$ is an even function, then the above relation leads to $e^{j\beta} = e^{-j\beta}$

implying that either $\beta = 0$ or $\beta = \pi$

- From $H(e^{j\omega}) = e^{j(c\omega+\beta)} \check{H}(\omega)$, we have $\check{H}(\omega) = e^{-j(c\omega+\beta)} H(e^{j\omega})$
- Substituting the value of β in the above we get $\check{H}(\omega) = \pm e^{-jc\omega} H(e^{j\omega}) = \pm \sum_n h[n]e^{-j\omega(c+n)}$
- Replacing ω with in the previous equation we get $\check{H}(-\omega) = \pm \sum_{l} h[l] e^{j\omega(c+l)}$
- Let I = N n, we rewrite the above equation as $\check{H}(-\omega) = \pm \sum_{n} h[N-n]e^{j\omega(c+N-n)}$
- As $\check{H}(\omega) = \check{H}(-\omega)$, we have $h[n]e^{-j\omega(c+n)} = h[N-n]e^{j\omega(c+N-n)}$
- The above leads to the condition:

 $h[n] = h[N - n], 0 \le n \le N$, where c = -N/2

 Thus, the FIR filter with an even amplitude response will <u>Original Pave at inear phase if it</u> has asymmetric impulse response-34

• If $\check{H}(\omega)$ is an odd function of ω , then from $e^{j(c\omega+\beta)}\check{H}(\omega) = e^{-j(c-\omega+\beta)}\check{H}(-\omega)$

we get $e^{j\beta} = -e^{-j\beta}$ as $\check{H}(-\omega) = -\check{H}(\omega)$

- The above is satisfied if $\beta = \pi/2$ or $\beta = -\pi/2$
- Then $H(e^{j\omega}) = e^{j(c\omega+\beta)} \check{H}(\omega)$ reduces to $H(e^{j\omega}) = je^{jc\omega} \check{H}(\omega)$
- The last equation can be rewritten as

$$\dot{\mathcal{H}}(\omega) = -je^{-jc\omega} H(e^{j\omega}) = -j \sum_n h[n]e^{-j\omega(c+n)}$$

- As $\check{H}(-\omega) = -\check{H}(\omega)$, from the above we get $\check{H}(-\omega) = j \sum_{l} h[l] e^{j\omega(c+l)}$
- Making a change of variable I = N n we have $\check{H}(-\omega) = j \sum_{n} h[N-n]e^{j\omega(c+N-n)}$

• Equating the above with

$$\check{H}(\omega) = -j \sum_{l} h[l] e^{-j\omega(c+l)}$$

we arrive at the condition for linear phase as

 $h[n] = h[N - n], 0 \le n \le N$, with c = -N/2

- Therefore, an FIR filter with an odd amplitude response will have linear-phase response if it has an antisymmetric impulse response
- Since the length of the impulse response can be either even or odd, we can define four types of linear-phase FIR transfer functions
- For an antisymmetric FIR filter of odd length, i.e., *N* even

$$h[N/2] = 0$$

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Linear-Phase FIR Transfer Functions



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Type 1: Symmetric Impulse Response with Odd Length

- In this case, the degree N is even (Assume N = 8)
- The transfer function H(z) is given by $H(z) = h[0] + h[1]z^{-1} + h[2] z^{-2} + h[3] z^{-3} + h[4] z^{-4} + h[5]z^{-5} + h[6]z^{-6} + h[7]z^{-7} + h[8]z^{-8}$
- Because of symmetry, we have h[0] = h[8], h[1] = h[7], h[2] = h[6], and h[3] = h[5]
- Thus, we can write

$$\begin{split} H(z) &= h[0](1+z^{-8}) + h[1](z^{-1}+z^{-7}) + h[2](z^{-2}+z^{-6}) + h[3](z^{-3}+z^{-5}) + h[4] z^{-4} \end{split}$$

$$= z^{-4} \{ h[0](z^4 + z^{-4}) + h[1](z^3 + z^{-3}) + h[2](z^2 + z^{-2}) + h[3](z + z^{-1}) + h[4] \}$$

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- The corresponding frequency response is then given by $H(e^{j\omega}) = e^{-j4\omega} \{2h[0]\cos(4\omega) + 2h[1]\cos(3\omega) + 2h[2]\cos(2\omega) + 2h[3]\cos(\omega) + h[4]\}$
- The quantity inside the braces is a real function of ω, and can assume positive or negative values in the range 0 ≤ |ω| ≤ π
- The phase function here is given by

 $\theta(\omega) = -4\omega + \beta$

where β is either 0 or $\pi,$ and hence it is a linear function of ω

• The group delay is given by

$$t(\omega) = -\frac{d\theta(\omega)}{d\omega} = 4$$

indicating a constant group delay of 4 samples

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• In the general case for Type 1 FIR filters, the frequency response is of the form

$$H(e^{j\omega}) = e^{-jN\omega/2}\tilde{H}(\omega)$$

where the **amplitude response** $\tilde{H}(\omega)$, also called the **zerophase response**, is of the form

$$\widetilde{H}(\omega) = h[\frac{N}{2}] + 2\sum_{n=1}^{N/2} h[\frac{N}{2} - n]\cos(\omega n)$$

• Example – Consider

$$H_0(z) = \frac{1}{6} \left[\frac{1}{2} + z^{-1} + z^{-2} + z^{-3} + z^{-4} + z^{-5} + \frac{1}{2} z^{-6} \right]$$

which is seen to be a slightly modified version of a length-7 moving-average FIR filter

• The above transfer function has a symmetric impulse response and therefore a linear phase response Original PowerPoint slides prepared by S. K. Mitra

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• The magnitude response of $H_0(z)$:



- Improved magnitude response obtained by changing the first and the last impulse response coefficients of MA filter
- This filter can be expressed as a cascade of a 2-point MA filter with a 6-point MA filter

$$H_0(z) = \frac{1}{2}(1+z^{-1}) \cdot \frac{1}{6}(1+z^{-1}+z^{-2}+z^{-3}+z^{-4}+z^{-5})$$

• Thus, $H_0(z)$ has a double zero at z = -1, i.e., ($\omega = \pi$) Original PowerPoint slides prepared by S. K. Mitra

Type 2: Symmetric Impulse Response with Even Length

- In this case, the degree N is odd (Assume N = 7)
- The transfer function H(z) is of the form

 $H(z) = h[0] + h[1]z^{-1} + \dots + h[7]z^{-7}$

- Because of symmetry, we have h[0] = h[7], h[1] = h[6], h[2] = h[5], and h[3] = h[4]
- Thus, we can write

$$\begin{split} H(z) &= h[0](1+z^{-7}) + h[1](z^{-1}+z^{-6}) + h[2](z^{-2}+z^{-5}) + h[3](z^{-3}\\&+z^{-4}) \\ &= z^{-7/2} \{h[0](z^{7/2}+z^{-7/2}) + h[1](z^{5/2}+z^{-5/2}) + h[2](z^{3/2}+z^{-3/2})\\&+ h[3](z^{1/2}+z^{-1/2}) \} \end{split}$$

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- The corresponding frequency response is then given by $H(e^{j\omega}) = e^{-j7\omega/2} \{2h[0]\cos(7\omega/2) + 2h[1]\cos(5\omega/2) + 2h[2]\cos(3\omega/2) + 2h[3]\cos(\omega/2)\}$
- The quantity inside the braces is a real function of ω, and can assume positive or negative values in the range 0 ≤ |ω| ≤ π
- The phase function here is given by

 $\theta(\omega) = -7\omega/2 + \beta$

where β is either 0 or $\pi,$ and hence it is a linear function of ω

• The group delay is given by

$$\tau(\omega) = \frac{7}{2}$$

indicating a constant group delay of 7/2 samples Original PowerPoint slides prepared by S. K. Mitra

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The expression for the frequency response in the general case for Type 2 FIR filters is of the form

 $H(e^{j\omega}) = e^{-jN\omega/2}\tilde{H}(\omega)$

• where the amplitude response is given by

$$\widetilde{H}(\omega) = 2 \sum_{n=1}^{(N+1)/2} h[\frac{N+1}{2} - n] \cos(\omega(n - \frac{1}{2}))$$

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Type 3: Antiymmetric Impulse Response with Odd Length

- In this case, the degree N is even (Assume N = 8)
- The transfer function H(z) is of the form

 $H(z) = h[0] + h[1]z^{-1} + \dots + h[8]z^{-8}$

- Antisymmetric filter coefficients: *h*[0] = −*h*[8], *h*[1] = −*h*[7], *h*[2] = −*h*[6], *h*[3] = − *h*[5], and *h*[4] = 0
- Applying the symmetry condition we get $H(z) = z^{-4} \{h[0](z^4 - z^{-4}) + h[1](z^3 - z^{-3}) + h[2](z^2 - z^{-2}) + h[3](z - z^{-1})\}$
- The corresponding frequency response is given by

 $H(e^{j\omega}) = e^{-j4\omega}e^{j\pi/2} \{2h[0]\sin(4\omega) + 2h[1]\sin(3\omega) + 2h[2]\sin(2\omega) + 2h[3]\sin(\omega)\}$

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• It also exhibits a linear phase response given by

$$\theta(\omega) = -4\omega + \pi/2 + \beta$$

where β is either 0 or π

• The group delay here is

$$T(\omega) = 4$$

indicating a constant group delay of 4 samples

• In the general case

$$H(e^{j\omega}) = je^{-jN\omega/2}\tilde{H}(\omega)$$

where the amplitude response is of the form

$$\widetilde{H}(\omega) = 2\sum_{n=1}^{N/2} h[\frac{N}{2} - n]\sin(\omega n)$$

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Type 4: Antiymmetric Impulse Response with Even Length

- In this case, the degree N is Odd (Assume N = 7)
- The transfer function H(z) is of the form

 $H(z) = h[0] + h[1]z^{-1} + \dots + h[7]z^{-7}$

- Antisymmetric filter coefficients: *h*[0] = −*h*[7], *h*[1] = −*h*[6], *h*[2] = −*h*[5], and *h*[3] = − *h*[4]
- Applying the symmetry condition we get $H(z) = z^{-7/2} \{h[0](z^{7/2} - z^{-7/2}) + h[1](z^{5/2} - z^{-5/2}) + h[2](z^{3/2} - z^{-3/2}) + h[3](z^{1/2} - z^{-1/2})\}$
- The corresponding frequency response is given by $H(e^{j\omega}) = e^{-j7\omega/2}e^{j\pi/2} \{2h[0]\sin(7\omega/2) + 2h[1]\sin(5\omega/2) + 2h[2]\sin(3\omega/2) + 2h[3]\sin(\omega/2)\}$

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General Form of Frequency Response

• In each of the four types of linear-phase FIR filters, the frequency response is of the form $II(-i\Theta) = -\frac{iN\Theta}{2} - \frac{i\Theta}{i\Theta} \tilde{I}(-)$

 $H(e^{j\omega}) = e^{-jN\omega/2} e^{j\beta} \tilde{H}(\omega)$

- The amplitude response $\tilde{H}(\omega)$ for each of the four types of linear-phase FIR filters can become negative over certain frequency ranges, typically in the stopband
- The magnitude and phase responses of the linear-phase FIR are given by

$$(e^{j\omega})| = |\tilde{H}(\omega)|$$

$$\theta(\omega) = \begin{cases} -\frac{N\omega}{2} + \beta, & \text{for } \tilde{H}(\omega) \ge 0\\ -\frac{N\omega}{2} + \beta - \pi, & \text{for } \tilde{H}(\omega) < 0 \end{cases}$$

• The group delay in each case is $\tau(\omega) = N/2$

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|H|

General Form of Frequency Response

- Note that, even though the group delay is constant, since in general |*H*(*e^{jω}*)| is not a constant, the output waveform is not a replica of the input waveform
- An FIR filter with a frequency response that is a real function of ω is often called a zero-phase filter
- Such a filter must have a noncausal impulse response

- Consider first an FIR filter with a symmetric impulse response: h[n] = h[N - n]
- Its transfer function can be written as

$$H(z) = \sum_{n=0}^{N} h[n] z^{-n} = \sum_{n=0}^{N} h[N-n] z^{-n}$$

• By making a change of variable m = N - n, we can write

$$\sum_{n=0}^{N} h[N-n]z^{-n} = \sum_{m=0}^{N} h[m]z^{-N+m} = z^{-N} \sum_{m=0}^{N} h[m]z^{m}$$

• But,

$$\sum_{m=0}^{N} h[m] z^{m} = H(z^{-1})$$

• Hence for an FIR filter with a symmetric impulse response of length *N*+1 we have $H(z) = z^{-N}H(z^{-1})$

• Such kind of *H*(*z*) is called a **mirror-image polynomial (MIP)** Original PowerPoint slides prepared by S. K. Mitra

- Now consider an FIR filter with a antisymmetric impulse response: h[n] = -h[N - n]
- Its transfer function can be written as

$$H(z) = \sum_{n=0}^{N} h[n] z^{-n} = -\sum_{n=0}^{N} h[N-n] z^{-n}$$

• By making a change of variable m = N - n, we get

$$-\sum_{n=0}^{N} h[N-n]z^{-n} = -\sum_{m=0}^{N} h[m]z^{-N+m} = -z^{-N}H(z^{-1})$$

- Hence, the transfer function H(z) of an FIR filter with an antisymmetric impulse response satisfies the condition $H(z) = -z^{-N}H(z^{-1})$
- A real-coefficient polynomial H(z) satisfying the above condition is called an antimirror-image polynomial (AIP)

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- It follows from the relation $H(z) = \pm z^{-N}H(z^{-1})$ that if $z = \xi_0$ is a zero of H(z), so is $z = 1/\xi_0$
- Moreover, for an FIR filter with a real impulse response, the zeros of H(z) occur in complex conjugate pairs
- Hence, a zero at $z = \xi_0$ is associated with a zero at $z = \xi_0^*$
- Thus, a complex zero that is not on the unit circle is ۲ associated with a set of 4 zeros given by

 $z = re^{\pm j\phi}, z = (1/r)e^{\pm j\phi}$

- A zero on the unit circle appear as a pair $z = e^{\pm j\phi}$, as its reciprocal is also its complex conjugate
- Since a zero at $z = \pm 1$ is its own reciprocal, it can appear only singly 7-52

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- Now a Type 2 FIR filter satisfies $H(z) = z^{-N}H(z^{-1})$ with degree *N* odd
- Hence, *H*(−1) = (−1)^{-N} *H*(−1) = −*H*(−1) implying that *H*(−1) = 0, i.e., *H*(*z*) must have a zero at *z* = −1
- a Type 3 or 4 FIR filter satisfies

$$H(z) = -z^{-N}H(z^{-1})$$

- Thus, $H(1) = -(1)^{-N} H(1) = -H(1)$ implying that H(z) must have a zero at z = 1
- On the other hand, only the Type 3 FIR filter is restricted to have a zero at z = −1 since here the degree N is even and hence, H(−1) = −(−1)^{-N} H(−1) = −H(−1)

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• Typical zero locations shown below



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Summary

- 1. Type 1 FIR filter: Either an even number or no zeros at z = 1 and z = -1
- 2. Type 2 FIR filter: Either an even number or no zeros at z = 1, and an odd number of zeros at z = -1
- 3. Type 3 FIR filter: An odd number of zeros at z = 1 and z = -1
- 4. Type 4 FIR filter: An odd number of zeros at z = 1, and either an even number or no zeros at z = -1

- The presence of zeros at $z = \pm 1$ leads to the following limitations on the use of these linear-phase transfer functions for designing frequency-selective filters
- A Type 2 FIR filter cannot be used to design a highpass filter since it always has a zero z = −1
- A Type 3 FIR filter has zeros at both at both z = 1 and z = -1, and hence cannot be used to design either a lowpass or a highpass or a bandstop filter
- A Type 4 FIR filter is not appropriate to design lowpass and bandstop filters due to the presence of a zero at z = 1
- Type 1 FIR filter has no such restrictions and can be used to design almost any type of filters

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Simple Lowpass FIR Digital Filters

 The simplest lowpass FIR digital filter is the 2-point movingaverage filter given by

$$H_0(z) = \frac{1}{2}(1+z^{-1}) = \frac{z+1}{2z}$$

- The above transfer function has a zero at z = -1 and a pole at z = 0
- Note that here the pole vector has a unity magnitude for all values of $\boldsymbol{\omega}$
- On the other hand, as ω increases from 0 to π, the magnitude of the zero vector decreases from a value of 2, the diameter of the unit circle, to 0
- Hence, the magnitude response $|H_0(e^{j\omega})|$ is a monotonically decreasing function of ω from $\omega = 0$ to $\omega = \pi$

- The maximum value of the magnitude function is 1 at $\omega = 0$, and the minimum value is 0 at $\omega = \pi$, i.e., $|H_0(e^{j0})| = 1$, $|H_0(e^{j\pi})| = 0$
- The frequency response of the above filter is given by $|H_0(e^{j\omega})| = e^{-j\omega/2}\cos(\omega/2)$
- The magnitude response $|H_0(e^{j\omega})| = e^{-j\omega/2} \cos(\omega/2)$ is a monotonically decreasing function of ω



• The frequency $\omega = \omega_c$ at which

$$|H_0(e^{j\omega_c})| = \frac{1}{\sqrt{2}} |H_0(e^{j0})|$$

is of practical interest since here the gain $G(\omega_c)$ in dB is given by $G(\omega_c) = 20\log_{10} |H(e^{j\omega_c})|$ $= 20\log_{10} |H(e^{j0})| - 20\log_{10} \sqrt{2} \approx -3 \text{ dB}$

since the dc gain $G(0) = 20\log_{10}|H_0(e^{j0})| = 0$

- Thus, the gain $G(\omega)$ at $\omega = \omega_c$ is approximately 3 dB less than the gain at $\omega = 0$
- As a result, ω_c is called the **3-dB cutoff frequency**
- To determine the value of ω_c we set $|H_0(e^{j\omega_c})|^2 = \cos^2(\omega_c/2) = \frac{1}{2}$ which yields $\omega_c = \pi/2$

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- The 3-dB cutoff frequency ω_{c} can be considered as the passband edge frequency
- As a result, for the filter $H_0(z)$ the passband width is approximately $\pi/2$, and the stopband is from $\pi/2$ to π
- Note: $H_0(z)$ has a zero at z = -1 or $\omega = \pi$, which is in the stopband of the filter
- A cascade of 3 sections of the FIR filter $H_0(z) = \frac{1}{2}(1 + z^{-1})$ results in an improved lowpass frequency



Highpass FIR Digital Filters

- The simplest highpass FIR filter is obtained from the simplest lowpass FIR filter by replacing *z* with -zFor example: $H_1(z) = \frac{1}{2}(1 - z^{-1})$
- Corresponding frequency response is given by
- $H_1(e^{j\omega}) = je^{-j\omega/2} \sin(\omega/2)$



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Highpass FIR Digital Filters

The monotonically increasing behavior of the magnitude function can again be demonstrated by examining the pole-zero pattern of the transfer function H₁(z)
For example: H₁(z) = ½(1 - z⁻¹)

• The highpass transfer function $H_1(z)$ has a zero at z = 1 or $\omega = 0$ which is in the stopband of the filter

- Improved highpass magnitude response can be obtained by cascading several sections of the first-order highpass filter
- Alternately, a higher-order highpass filter of the form $H_1(z) = \frac{1}{M} \sum_{n=0}^{M-1} (-1)^n z^{-n}$

is obtained by replacing z with -z in the transfer function of a moving average filter

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Highpass FIR Digital Filters

- An application of the FIR highpass filters is in moving-targetindicator (MTI) radars
- In these radars, interfering signals, called **clutters**, are generated from fixed objects in the path of the radar beam
- The clutter, generated mainly from ground echoes and weather returns, has frequency components near zero frequency (dc)
- The clutter can be removed by filtering the radar return signal through a **two-pulse canceler**, which is the first-order FIR highpass filter $H_1(z) = \frac{1}{2}(1 z^{-1})$
- For a more effective removal it may be necessary to use a three-pulse canceler obtained by cascading two two-pulse cancelers

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Simple IIR Digital Filters

• We have shown earlier that the first-order causal IIR transfer function $H(z) = \frac{K}{1 - \alpha z^{-1}}, \quad 0 < \alpha < 1$

has a lowpass magnitude response for $\alpha > 0$

• On the other hand, the first-order causal IIR transfer function $H(z) = \frac{K}{1 - \alpha z^{-1}}, -1 < \alpha < 0$

has a highpass magnitude response for $\alpha < 0$

• However, the modified transfer function obtained with the addition of a factor $(1 + z^{-1})$ to the numerator

$$H(z) = \frac{K(1 + z^{-1})}{1 - \alpha z^{-1}}, \quad -1 < \alpha < 0$$

exhibits a lowpass magnitude response

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Simple IIR Digital Filters

• The modified first-order lowpass transfer function for both positive and negative values of α is then given by

$$H_{LP}(z) = \frac{K(1+z^{-1})}{1-\alpha z^{-1}}, \quad 0 < |\alpha| < 1$$

- As ω increases from 0 to $\pi,$ the magnitude of the zero vector decreases from a value of 2 to 0
- The maximum values of the magnitude function is $2K/(1-\alpha)$ at $\omega = 0$ and the minimum value is 0 at $\omega = \pi$, i.e.,

 $|H_{LP}(e^{j0})| = 2K/(1-\alpha), |H_{LP}(e^{j\pi})| = 0$

• Therefore, $|H_{LP}(e^{j\omega})|$ is a monotonically decreasing function of ω from $\omega = 0$ to $\omega = \pi$

Simple IIR Digital Filters

- For most applications, it is usual to have a dc gain of 0 dB, that is to have $|H_{LP}(e^{j0})| = 1$
- To this end, we choose $K = (1 \alpha)/2$ resulting in the firstorder IIR lowpass transfer function

$$H_{LP}(z) = \frac{1 - \alpha}{2} \left(\frac{1 + z^{-1}}{1 - \alpha z^{-1}} \right), \quad 0 < |\alpha| < 1$$

• The above transfer function has a zero at i.e., at $\omega = \pi$ which is in the stopband

• A first-order causal lowpass IIR digital filter has a transfer function given by $1-\alpha(1+z^{-1})$

$$H_{LP}(z) = \frac{1 - \alpha}{2} \left(\frac{1 + z^{-1}}{1 - \alpha z^{-1}} \right)$$

where $|\alpha| < 1$ for stability

- The above transfer function has a zero at z = -1 i.e., at ω = π which is in the stopband
- $H_{LP}(z)$ has a real pole at $z = \alpha$
- As ω increases from 0 to π, the magnitude of the zero vector decreases from a value of 2 to 0, whereas, for a positive value of α, the magnitude of the pole vector increases from a value of 1- α to 1+ α function is 1 at ω = 0, and the minimum
- The maximum value of the magnitude function is 1 at $\omega = 0$, and the minimum value is 0 at $\omega = \pi$ Original PowerPoint slides prepared by S. K. Mitra

- That is $|H_{LP}(e^{j0})| = 1$, $|H_{LP}(e^{j\pi})| = 0$
- Therefore, $|H_{LP}(e^{j\omega})|$ is a monotonically decreasing function of ω from $\omega = 0$ to $\omega = \pi$ as indicated below



• The squared magnitude function is given by

$$|H_{LP}(e^{j\omega})|^{2} = \frac{(1-\alpha)^{2}(1+\cos\omega)}{2(1+\alpha^{2}-2\alpha\cos\omega)}$$

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• The derivative of $|H_{LP}(e^{j\omega})|^2$ with respect to ω is given by

$$\frac{d |H_{LP}(e^{j\omega})|^2}{d\omega} = \frac{-(1-\alpha)^2(1+2\alpha+\alpha^2)\sin\omega}{2(1-2\alpha\cos\omega+\alpha^2)^2}$$

 $d|H_{LP}(e^{j\omega})|^2 / d\omega \le 0$ in the range $0 \le \omega \le \pi$ verifying again the monotonically decreasing behavior of the magnitude function

• To determine the 3-dB cutoff frequency we set

$$|H_{LP}(e^{j\omega_c})|^2 = \frac{1}{2}$$

• in the expression for the square magnitude function resulting in $\frac{(1-\alpha)^2(1+\cos\omega_c)}{2} = \frac{1}{2}$

$$\frac{1}{2(1+\alpha^2-2\alpha\cos\omega_c)}=\frac{1}{2}$$

or
$$(1 - \alpha)^2 (1 + \cos \omega_c) = 1 + \alpha^2 - 2\alpha \cos \omega_c \Rightarrow \cos \omega_c = 2\alpha/(1 + \alpha^2)$$

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- The above quadratic equation can be solved for α yielding two solutions
- The solution resulting in a stable transfer function $H_{LP}(z)$ is given by $\alpha = \frac{1 \sin \omega_c}{\cos \omega_c}$
- It follows from

$$H_{LP}(e^{j\omega})|^{2} = \frac{(1-\alpha)^{2}(1+\cos\omega)}{2(1+\alpha^{2}-2\alpha\cos\omega)}$$

that $H_{LP}(z)$ is a BR function for $|\alpha| < 1$

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Highpass IIR Digital Filters

• A first-order causal highpass IIR digital filter has a transfer function given by

$$H_{HP}(z) = \frac{1+\alpha}{2} \left(\frac{1-z^{-1}}{1-\alpha z^{-1}} \right)$$

where $|\alpha| < 1$ for stability

- The above transfer function has a zero at z = 1 i.e., at ω = 0 which is in the stopband
- Its 3-dB cutoff frequency is given by

 $\alpha = (1 - \sin \omega_c) / \cos \omega_c$

which is the same as that of $H_{LP}(z)$

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Highpass IIR Digital Filters

• Magnitude and gain responses of $H_{HP}(z)$ are shown below



• $H_{HP}(z)$ is a BR function for $|\alpha| < 1$

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Highpass IIR Digital Filters

- Magnitude and gain responses of $H_{HP}(z)$ are shown below
- Example Design a first-order highpass digital filter with a 3dB cutoff frequency of 0.8π
- Now, $\sin(\omega_c) = \sin(0.8\pi) = 0.587785$ and $\cos(0.8\pi) = -0.80902$
- Therefore $\alpha = (1 \sin \omega_c)/\cos \omega_c = -0.5095245$
- Therefore,

$$H_{HP}(z) = \frac{1+\alpha}{2} \left(\frac{1-z^{-1}}{1-\alpha z^{-1}} \right)$$
$$= 0.245238 \left(\frac{1-z^{-1}}{1+0.5095245 z^{-1}} \right)$$

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• A 2nd-order bandpass digital transfer function is given by

$$H_{BP}(z) = \frac{1 - \alpha}{2} \left(\frac{1 - z^{-2}}{1 - \beta(1 + \alpha) z^{-1} + \alpha z^{-2}} \right)$$

Its squared magnitude function is

$$= \frac{\left|H_{BP}(e^{j\omega})\right|^{2}}{2[1+\beta^{2}(1+\alpha)^{2}+\alpha^{2}-2\beta(1+\alpha)^{2}\cos\omega+2\alpha\cos2\omega]}$$

- $|H_{BP}(e^{j\omega})|^2$ goes to zero at $\omega = 0$ and $\omega = \pi$
- It assumes a maximum value of 1 at $\omega = \omega_o$ called the **center frequency** of the bandpass filter, where

$$\omega_o = \cos^{-1}(\beta)$$

• The frequencies ω_{c1} and ω_{c2} where $|H_{BP}(e^{j\omega})|^2$ becomes 1/2 Original Green called the 3-d B cutoff frequencies 7-74

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 The difference between the two cutoff frequencies, assuming ω_{c1} > ω_{c2} is called the **3-dB bandwidth** and is given by

$$B_w = \omega_{c2} - \omega_{c1} = \cos^{-1} \left(\frac{2\alpha}{1 + \alpha^2} \right)$$

• The transfer function $H_{BP}(z)$ is a BR function if $|\alpha| < 1$ and $|\beta| < 1$



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- Example Design a 2nd order bandpass digital filter with central frequency at 0.4π and a 3-dB bandwidth of 0.1π
- Here $\beta = \cos(\omega_o) = \cos(0.4\pi) = 0.309017$ and $2\alpha/(1+\alpha)^2 = \cos(B_w) = \cos(0.1\pi) = 0.9510565$
- The solution of the above equation yields: $\alpha = 1.376382$ and $\alpha = 0.72654253$
- The corresponding transfer functions are

$$H'_{BP}(z) = -0.18819 \frac{1 - z^{-2}}{1 - 0.7343424z^{-1} + 1.37638z^{-2}}$$

and
$$H''_{BP}(z) = 0.13673 \frac{1 - z^{-2}}{1 - 0.533531z^{-1} + 0.72654253z^{-2}}$$

• The poles of $H'_{BP}(z)$ are at $z = 0.3671712 \pm j1.11425636$ and have a magnitude > 1

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- Thus, the poles of $H_{BP}(z)$ are outside the unit circle making the transfer function unstable
- On the other hand, the poles of $H_{BP}^{"}(z)$ are at $z = 0.2667655 \pm j0.8095546$ and have a
- The solution of the above equation yields: $\alpha = 1.376382$ and $\alpha = 0.72654253$ and have a magnitude of 0.8523746
- Hence $H_{BP}^{"}(z)$ is BIBO stable



Bandstop IIR Digital Filters

A 2nd-order bandstop digital filter has a transfer function given by

$$H_{BS}(z) = \frac{1+\alpha}{2} \left(\frac{1-2\beta z^{-1} + z^{-2}}{1-\beta(1+\alpha) z^{-1} + \alpha z^{-2}} \right)$$

• The transfer function $H_{BS}(z)$ is a BR function if $|\alpha| < 1$ and $|\beta| < 1$



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Bandstop IIR Digital Filters

- Here, the magnitude function takes the maximum value of 1 at ω = 0 and ω = π
- It goes to 0 at $\omega = \omega_o$, where ω_o , called the **notch frequency**, is given by

 $\omega_o = \cos(-1)(\beta)$

- The digital transfer function $H_{BS}(z)$ is more commonly called a **notch filter**
- The frequencies ω_{c1} and ω_{c2} where |H_{BS}(e^{jω})|² becomes 1/2 are called the 3-dB cutoff frequencies
- The difference between the two cutoff frequencies, assuming $\omega_{c1} > \omega_{c2}$ is called the **3-dB notch bandwidth** and is given by

$$B_w = \omega_{c2} - \omega_{c1} = \cos^{-1} \left(\frac{2\alpha}{1 + \alpha^2} \right)$$

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- By cascading the simple digital filters discussed so far, we can implement digital filters with sharper magnitude responses
- Consider a cascade of *K* first-order lowpass sections characterized by the transfer function

$$H_{LP}(z) = \frac{1 - \alpha}{2} \left(\frac{1 + z^{-1}}{1 - \alpha z^{-1}} \right)$$

• The corresponding squared-magnitude function is given by

$$|G_{LP}(e^{j\omega})|^{2} = \left[\frac{(1-\alpha)^{2}(1+\cos\omega)}{2(1+\alpha^{2}-2\alpha\cos\omega)}\right]^{K}$$

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• To determine the relation between its 3-dB cutoff frequency ω_c and the parameter α , we set

$$\left[\frac{(1-\alpha)^2(1+\cos\omega_c)}{2(1+\alpha^2-2\alpha\cos\omega_c)}\right]^K = \frac{1}{2}$$

which when solved for α , yields for a stable $G_{LP}(z)$

$$\alpha = \frac{1 + (1 - C)\cos\omega_c - \sin\omega_c\sqrt{2C - C^2}}{1 - C + \cos\omega_c}$$

where $C = 2^{(K-1)/K}$

• It should be noted that the expression for α given earlier reduces to $1-\sin \omega$.

$$\alpha = \frac{1 - \sin \omega_c}{\cos \omega_c}$$

for
$$K = 1$$

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- Example Design a lowpass filter with a 3- dB cutoff frequency at $\omega_c = 0.4\pi$ using a single first-order section and a cascade of 4 first-order sections, and compare their gain responses
- For the single first-order lowpass filter we have

$$\alpha = \frac{1 + \sin \omega_c}{\cos \omega_c} = \frac{1 + \sin(0.4\pi)}{\cos(0.4\pi)} = 0.1584$$

 For the cascade of 4 first-order sections, we substitute K = 4 and get

$$C = 2^{(K-1)/K} = 2^{(4-1)/4} = 1.6818$$

• Next we compute $\alpha = \frac{1 + (1 - C)\cos\omega_c - \sin\omega_c\sqrt{2C - C^2}}{1 - C + \cos\omega_c}$ Original PowerPoint slides prepared by S. The McGraw-Hill Companies = -0.251

- The gain responses of the two filters are shown below
- As can be seen, cascading has resulted in a sharper roll-off in the gain response



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- The simple filters discussed so far are characterized either by a single passband and/or a single stopband
- There are applications where filters with multiple passbands and stopbands are required
- The **comb filter** is an example of such filters
- In its most general form, a comb filter has a frequency response that is a periodic function of ω with a period $2\pi/L$, where *L* is a positive integer
- If H(z) is a filter with a single passband and/or a single stopband, a comb filter can be easily generated from it by replacing each delay in its realization with *L* delays resulting in a structure with a transfer function given by $G(z) = H(z^L)$

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- If $|H(e^{j\omega})|$ exhibits a peak at ω_p , then $|G(e^{j\omega})|$ will exhibit *L* peaks at $\omega_p k/L$, $0 \le k \le L 1$ in the frequency range $0 \le \omega < 2\pi$
- Likewise, if $|H(e^{j\omega})|$ has a notch at ω_o , then then $|G(e^{j\omega})|$ have L notches at $\omega_o k/L$, $0 \le k \le L - 1$ in the frequency range $0 \le \omega \le 2\pi$
- A comb filter can be generated from either an FIR or an IIR prototype filter
- For example, the comb filter generated from the prototype lowpass FIR filter $H_0(z) = 1/2 (1 + z^{-1})$ has a transfer function $G_0(z) = H_0(z^L) = 1/2 (1 + z^{-L})$
- $|G_0(e^{j\omega})|$ has *L* notches at $\omega = (2k+1)\pi/L$ and *L* peaks at $\omega = 2\pi k/L$, $0 \le k \le L-1$, in the frequency range $0 \le \omega < 2\pi$ Original PowerPoint slides prepared by S. K. Mitra

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• On the other hand, the comb filter generated from the prototype lowpass FIR filter $H_1(z) = 1/2 (1 - z^{-1})$ has a transfer function

$$G_1(z) = H_1(z^L) = 1/2 (1 - z^{-L})$$

• $|G_1(e^{j\omega})|$ has L notches at $\omega = (2k+1)\pi/L$ and L peaks at $\omega = 2\pi k/L$, $0 \le k \le L-1$, in the frequency range $0 \le \omega < 2\pi$



- Depending on applications, comb filters with other types of periodic magnitude responses can be easily generated by appropriately choosing the prototype filter
- For example, the *M*-point moving average filter

$$H(z) = \frac{1 - z^{-M}}{M(1 - z^{-1})}$$

has been used as a prototype

- This filter has a peak magnitude at $\omega = 0$, and M 1 notches at $\omega = 2\pi I/M$, $1 \le I \le M 1$
- The corresponding comb filter has a transfer function

$$G(z) = \frac{1 - z^{-LM}}{M(1 - z^{-L})}$$

whose magnitude has L peaks at $2\pi k/L$, $1 \le k \le L - 1$ and L(M - 1) notches at $2\pi k/LM$, $1 \le k \le L(M - 1)$

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 A set of *L* transfer functions, {*H_i(z)*}, 0 ≤ *i* ≤ *L* − 1, is defined to be **delay-complementary** of each other if the sum of their transfer functions is equal to some integer multiple of unit delays, i.e., ^{L-1}
 ∑*H_i(z) = βz^{-n_o}*, β ≠ 0

where n_o is a nonnegative integer

i=0

- A delay-complementary pair {H₀(z), H₁(z)} can be readily designed if one of the pairs is a known Type 1 FIR transfer function of odd length
- Let $H_0(z)$ be a Type 1 FIR transfer function of length M = 2K+1, its delay-complementary transfer function is given by $H_1(z) = z^{-K} H_0(z)$

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- Let the magnitude response of $H_0(z)$ be equal to $1 \pm \delta_p$ in the passband and less than or equal to δ_s in the stopband where δ_p and δ_s are very small numbers
- Now the frequency response of $H_0(z)$ can be expressed as

$$H_0(e^{j\omega}) = e^{-jK\omega} \breve{H}_0(\omega)$$

where $H_0(\omega)$ is the **amplitude response**

• Its delay-complementary transfer function $H_1(z)$ has a frequency response given by

$$H_1(e^{j\omega}) = e^{-jK\omega} \breve{H}_1(\omega) = e^{-jK\omega} [1 - \breve{H}_0(\omega)]$$

• Now, in the passband, $1 - \delta_p \le H_0(\omega) \le 1 + \delta_p$, and in the stopband, $-\delta_s \le H_0(\omega) \le \delta_s$

- It follows from the above equation that $-\delta_p \leq H_1(\omega) \leq \delta_p$, and in the stopband, $1 \delta_s \leq H_1(\omega) \leq 1 + \delta_s$
- As a result, H₁(z) has a complementary magnitude response characteristic to that of H₀(z) with a stopband exactly identical to the passband of H₀(z), and a passband that is exactly identical to the stopband of H₀(z)
- Thus, if $H_0(z)$ is a lowpass filter, $H_1(z)$ will be a highpass filter, and vice versa
- The frequency ω_o at which $H_0(\omega_o) = H_1(\omega_o) = 0.5$ the gain responses of both filters are 6 dB below their maximum values
- The frequency is thus called the **6-dB crossover frequency**

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- Example Consider the Type 1 bandstop transfer function $H_{BS}(z) = \frac{1}{64}(1+z^{-2})^4(1-4z^{-2}+5z^{-4}+5z^{-8}-4z^{-10}+z^{-12})$
- Its delay-complementary Type 1 bandpass transfer function is given by $H_{BP}(z) = z^{-10} H_{BS}(z)$



 $=\frac{1}{64}(1-z^{-2})^4(1+4z^{-2}+5z^{-4}+5z^{-8}+4z^{-10}+z^{-12})$

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Allpass Complementary Filters

- A set of *M* transfer functions, $\{H_i(z)\}, 0 \le i \le M 1$, is defined to be **allpass-complementary** of each other, if the sum of their transfer functions is equal to an allpass function, i.e., $\sum_{i=1}^{M-1} H_i(z) = A(z)$
- Example Consider the two transfer functions $H_0(z)$ and $H_1(z)$ given by

 $H_0(z) = \frac{1}{2} [A_0(z) + A_1(z)]$ $H_1(z) = \frac{1}{2} [A_0(z) - A_1(z)]$

where $A_0(z)$ and $A_1(z)$ are stable allpass transfer functions

- Note that $H_0(z) + H_1(z) = A_0(z)$
- Hence, $H_0(z)$ and $H_1(z)$ are allpass complementary

Power Complementary Filters

- A set of *M* transfer functions, $\{H_i(z)\}, 0 \le i \le M 1$, is defined to be **power-complementary** of each other, if the sum of their square-magnitude responses is equal to a constant *K* for all values of ω , i.e., $\frac{M-1}{\sum_{i=0}^{M-1} |H_i(e^{j\omega})|^2} = K$, for all ω
- By analytic continuation, the above property is equal to $\sum_{i=0}^{M-1} H_i(z)H_i(z^{-1}) = K, \quad \text{for all } \omega$

for real coefficient $H_0(z)$

 Usually, by scaling the transfer functions, the powercomplementary property is defined for K = 1

Power Complementary Filters

- For a pair of power-complementary transfer functions, $H_0(z)$ and $H_1(z)$, the frequency ω_o where $|H_0(e^{j\omega o})|^2 = |H_1(e^{j\omega o})|^2 = 0.5$, is called the **cross-over frequency**
- At this frequency the gain responses of both filters are 3-dB below their maximum values
- As a result, is called the **3-dB cross-over frequency**
- Example Consider the two transfer functions $H_0(z)$ and $H_1(z)$ given by

 $H_0(z) = \frac{1}{2} [A_0(z) + A_1(z)]$ $H_1(z) = \frac{1}{2} [A_0(z) - A_1(z)]$

where $A_0(z)$ and $A_1(z)$ are stable allpass transfer functions

• $H_0(z)$ and $H_1(z)$ are allpass and power complementary

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Doubly Complementary Filters

- A set of *M* transfer functions satisfying both the allpass complementary and the power complementary properties is known as a doubly-complementary set
- A pair of doubly-complementary IIR transfer functions, $H_0(z)$ and $H_1(z)$, with a sum of allpass decomposition can be simply realized as indicated below



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Doubly Complementary Filters

Example - The first-order lowpass transfer function

$$H_{LP}(z) = \frac{1-\alpha}{2} \left(\frac{1+z^{-1}}{1-\alpha z^{-1}} \right)$$

can be expressed as

$$H_{LP}(z) = \frac{1}{2} \left(1 + \frac{-\alpha + z^{-1}}{1 - \alpha z^{-1}} \right) = \frac{1}{2} \left[A_0(z) + A_1(z) \right]$$

where

$$A_0(z) = 1, \quad A_1(z) = \frac{-\alpha + z^{-1}}{1 - \alpha z^{-1}}$$

• Its power-complementary highpass transfer function is thus given by $H_{HP}(z) = \frac{1}{2} [A_0(z) - A_1(z)] = \frac{1}{2} \left(1 - \frac{-\alpha + z^{-1}}{1 - \alpha z^{-1}} \right)$ $= \frac{1 + \alpha}{2} \left(\frac{1 - z^{-1}}{1 - \alpha z^{-1}} \right)$

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Doubly Complementary Filters

- The above expression is precisely the firstorder highpass transfer function described earlier
- Figure below demonstrates the allpass complementary property and the power complementary property of can be expressed as $H_{IP}(z)$ and $H_{HP}(z)$



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Power Symmetric Filters

 A real-coefficient causal digital filter with a transfer function H(z) is said to be a power-symmetric filter if it satisfies the condition

$$H(z)H(z^{-1}) + H(-z)H(-z^{-1}) = K$$

where K > 0 is a constant

- It can be shown that the gain function $G(\omega)$ of a powersymmetric transfer function at $\omega = \pi$ is given by
- $10\log_{10}K 3 \, dB$
- If we define G(z) = H(-z), then it follows from the definition of the power-symmetric filter that H(z) and G(z) are power-complementary as

$$H(z)H(z^{-1}) + G(z)G(z^{-1}) = a \text{ constant}$$

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Conjugate Quadratic Filters

• If a power-symmetric filter has an FIR transfer function H(z)of order *N*, then the FIR digital filter with a transfer function $G(z) = z^{-N}H(-z^{-1})$

is called a **conjugate quadratic filter** of H(z) and vice-versa

- It follows from the definition that *G*(*z*) is also a power-symmetric causal filter quadratic filters
- It also can be seen that a pair of conjugate quadratic filters
 H(z) and G(z) are also power-complementary
- Example Let $H(z) = 1 2z^{-1} + 6z^{-2} + 3z^{-3}$

 $H(z)H(z^{-1}) + H(-z)H(-z^{-1}) = 100$

• H(z) is a power-symmetric transfer function