

## Homework Solutions #7

(Due date: 2011/04/25)

This problem set covers materials of Lesson 9. The full score is 40 points.

- 1) (15%) Fig. 1 shows a simplified model of  $pn$ -junction. A depletion region is formed during  $-d_p < x < d_n$ , where net charge densities  $-eN_A$  ( $\text{C/m}^3$ ) and  $eN_D$  ( $\text{C/m}^3$ ) exist for  $-d_p < x < 0$  and  $0 < x < d_n$ , respectively. Note that  $N_A d_p = N_D d_n$  holds to maintain the electric neutrality. Let the potential  $V(-d_p) = 0$ , find the potential distribution  $V(x)$  for  $-d_p < x < d_n$  by solving Poisson's equation (with suitable boundary conditions).

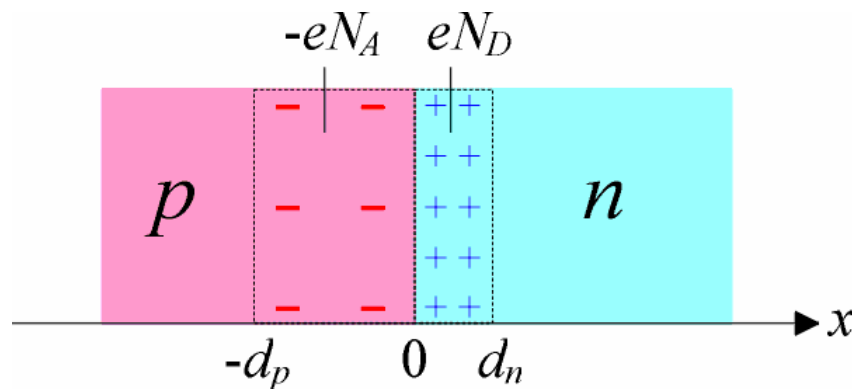


Fig. 1.  $pn$ -junction.

Answer:

We discuss the potential distribution during (i)  $-d_p < x < 0$  and (ii)  $0 < x < d_n$  separately, due to the discontinuous charge distribution at the interface.

(i)  $-d_p < x < 0$ :

According to the Poisson's equation, we have:

$$\frac{d^2 V_A}{dx^2} = \frac{eN_A}{\epsilon_A}, \Rightarrow V_A(x) = \frac{eN_A}{2\epsilon_A} x^2 + c_1 x + c_2.$$

Based on the BC:  $V_A(-d_p) = 0$ , we get  $\frac{eN_A}{2\epsilon_A} d_p^2 - c_1 d_p + c_2 = 0$ .

Another BC is that the potential is constant (E-field is zero) near  $x = -d_p$ , otherwise,

free electrons would keep on moving (not in steady state).

$$\Rightarrow V'_A(-d_p) = 0, \Rightarrow c_1 = \frac{d_p eN_A}{\epsilon_A}; \Rightarrow c_2 = \frac{eN_A}{2\epsilon_A} d_p^2.$$

So,  $V_A(x) = \frac{eN_A}{2\epsilon_A} x^2 + \frac{eN_A d_p}{\epsilon_A} x + \frac{eN_A}{2\epsilon_A} d_p^2$ , which is a quadratic function of  $x$ .

(ii)  $0 < x < d_n$ :

Similar with the procedure we just perform:

$$\frac{d^2 V_D}{dx^2} = -\frac{eN_D}{\epsilon_D}, \Rightarrow V_D(x) = -\frac{eN_D}{2\epsilon_D} x^2 + c_3 x + c_4.$$

And BC:  $V'_D(d_n) = 0$ ,  $\Rightarrow -\frac{eN_D}{\epsilon_D} d_n + c_3 = 0, \Rightarrow c_3 = \frac{eN_D}{\epsilon_D} d_n$ .

At the interface  $x = 0$ , the potential has to be continuous:

$$V_A(0) = V_D(0) \Rightarrow \frac{eN_A}{2\epsilon_A} d_p^2 = c_4.$$

So,  $V_D(x) = -\frac{eN_D}{2\epsilon_D} x^2 + \frac{eN_D d_n}{\epsilon_D} x + \frac{eN_A}{2\epsilon_A} d_p^2$

- 2) Consider a coaxial cable capacitor with the same geometry as that shown in Fig. 9-4 in the lecture notes, while the space between the conducting surfaces ( $a < r < b$ ) is filled with two dielectric media of permittivities  $\epsilon_1, \epsilon_2$  according to:

$$\epsilon = \begin{cases} \epsilon_1, & \text{for } \phi \in (0, \pi/2) \\ \epsilon_2, & \text{for } \phi \in (\pi/2, 2\pi) \end{cases} \quad (r, \phi \text{ are with cylindrical coordinates}).$$

- 2a) (10%) Deposit charges of  $+Q$  and  $-Q$  on the inner and outer conducting surfaces, respectively. Find the electric flux density  $\vec{D}$  in the dielectric region  $a < r < b$ .  
(Hint: Use boundary condition.)

Answer:

Let  $Q_1, Q_2$  represent the charges on the inner conducting surface in contact with Medium 1 and Medium 2, respectively. We also know that  $Q_1 + Q_2 = Q \dots (1)$ .

$$\oint \vec{D} \cdot d\vec{s} = \int_{V'} \rho_f dv' \Rightarrow \vec{D} = \begin{cases} \vec{a}_r \frac{2}{\pi L} Q_1, & \text{at medium 1 } (\epsilon = \epsilon_1) \\ \vec{a}_r \frac{2}{3\pi L} Q_2, & \text{at medium 2 } (\epsilon = \epsilon_2) \end{cases}.$$

$$\Rightarrow \vec{E} = \begin{cases} \vec{a}_r \frac{2}{\pi \epsilon_1 L} Q_1, & \text{at medium 1 } (\epsilon = \epsilon_1) \\ \vec{a}_r \frac{2}{3\pi \epsilon_2 L} Q_2, & \text{at medium 2 } (\epsilon = \epsilon_2) \end{cases}$$

By the boundary condition between Medium 1 and 2,  $E_{1t} = E_{2t}$ ,

$$\Rightarrow \frac{2}{\epsilon_1} Q_1 = \frac{2}{3\epsilon_2} Q_2 \dots (2)$$

With (1) and (2), we can derive that  $Q_1 = \frac{\epsilon_1}{\epsilon_1 + 3\epsilon_2} Q$ , and  $Q_2 = \frac{3\epsilon_2}{\epsilon_1 + 3\epsilon_2} Q$ .

$$\vec{D} = \begin{cases} \vec{a}_r \frac{2}{\pi L} \frac{\epsilon_1}{\epsilon_1 + 3\epsilon_2} Q, & \text{at medium 1 } (\epsilon = \epsilon_1) \\ \vec{a}_r \frac{2}{\pi L} \frac{\epsilon_2}{\epsilon_1 + 3\epsilon_2} Q, & \text{at medium 2 } (\epsilon = \epsilon_2) \end{cases}$$

- 2b) (5%) Find the corresponding capacitance  $C$ , and the effective permittivity  $\epsilon_{eff}$  if

$$\varepsilon_{eff} \equiv \frac{C \ln(b/a)}{2\pi L} \quad (\text{justified by the result of Example 9-3 in the lecture notes}).$$

Answer:

$$V = \int_+^- \vec{E} \cdot d\vec{l} = \int_a^b \vec{a}_r \frac{2}{\pi r L} \frac{1}{\varepsilon_1 + 3\varepsilon_2} Q \cdot \vec{a}_r dr = \frac{2Q}{\pi L(\varepsilon_1 + 3\varepsilon_2)} \int_a^b \frac{1}{r} dr = \frac{2Q}{\pi L(\varepsilon_1 + 3\varepsilon_2)} \ln\left(\frac{b}{a}\right),$$

$$C = \frac{Q}{V} = \frac{\pi L}{2} (\varepsilon_1 + 3\varepsilon_2) \frac{1}{\ln(b/a)},$$

$$\varepsilon_{eff} \equiv \frac{C \ln(b/a)}{2\pi L} = \frac{\pi L}{2} (\varepsilon_1 + 3\varepsilon_2) \frac{1}{\ln(b/a)} \frac{\ln(b/a)}{2\pi L} = \frac{\varepsilon_1 + 3\varepsilon_2}{4}.$$

3) (10%) Problem **P.3–32** of the textbook.

Answer:

Assumed that there are total  $+Q$  charge uniformly distributed on the surface of the inner metal. The positive charge will induce  $-Q$  deposit on the surface outer metal as shown by the figure.

For  $r_i < r < b$ :

From the Gauss law  $\oint \vec{E}_1 \cdot d\vec{s} = \vec{a}_r E_1 \cdot 2\pi r \cdot L = \frac{Q}{\epsilon_0 \epsilon_{r1}} \Rightarrow \vec{E}_1 = \vec{a}_r \frac{Q}{2\pi L \cdot \epsilon_0 \epsilon_{r1}}$ , where  $L$  is length of the transmission line.

$$V_1 = -\int_b^{r_i} \vec{E}_1 \cdot d\vec{r} = \vec{a}_r \frac{Q}{2\pi L \epsilon_0 \epsilon_{r1}} \ln\left(\frac{b}{r_i}\right),$$

$$\text{Then the capacitance per unit length is } \frac{C_1}{L} = \frac{Q}{V_1} = \frac{2\pi \epsilon_0 \epsilon_{r1}}{\ln\left(\frac{b}{r_i}\right)}.$$

For  $b < r < r_o$ :

$$\text{From the same procedure, the capacitance per unit length is } \frac{C_2}{L} = \frac{Q}{V_2} = \frac{2\pi \epsilon_0 \epsilon_{r2}}{\ln\left(\frac{r_o}{b}\right)}.$$

The total capacitance is the series connection of the two sub-capacitances, therefore

$$\frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_2}, \Rightarrow C = \frac{2\pi \epsilon_0 \epsilon_{r1} \epsilon_{r2}}{\epsilon_{r2} \ln\left(\frac{b}{r_i}\right) + \epsilon_{r1} \ln\left(\frac{r_o}{b}\right)}$$

